#### An Introduction to Symmetric Functions

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In other words, if  $i_1, \ldots, i_m$  are distinct positive integers and  $\alpha_1, \ldots, \alpha_m$  are arbitrary nonnegative integers then the coefficient of  $x_{i_1}^{\alpha_1} \cdots x_{i_m}^{\alpha_m}$  in a symmetric function is the same as the coefficient of  $x_1^{\alpha_1} \cdots x_m^{\alpha_m}$ .

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- $\blacktriangleright \sum_{i \leq j} x_i x_j$

But not  $\sum_{i \leq j} x_i x_j^2$ 

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- Symmetric functions are useful in counting unlabeled graphs (Pólya theory).

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There are several important bases for  $\Lambda^n$ , all indexed by partitions.

## Monomial symmetric functions

If a symmetric function has a term  $x_1^2 x_2 x_3$  with coefficient 1, then it must contain all terms of the form  $x_i^2 x_j x_k$ , with *i*, *j*, and *k* distinct, with coefficient 1. If we add up all of these terms, we get the monomial symmetric function

$$m_{(2,1,1)} = \sum x_i^2 x_j x_k$$

where the sum is over all distinct terms of the form  $x_i^2 x_j x_k$  with *i*, *j*, and *k* distinct.

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$$m_{(2,1,1)} = x_1^2 x_2 x_3 + x_3^2 x_1 x_4 + x_1^2 x_3 x_5 + \cdots$$

More generally, for any partition  $\lambda = (\lambda_1, ..., \lambda_k)$ ,  $m_{\lambda}$  is the sum of all distinct monomials of the form

 $x_{i_1}^{\lambda_1}\cdots x_{i_k}^{\lambda_k}.$ 

It's easy to see that  $\{m_{\lambda}\}_{\lambda \vdash n}$  is a basis for  $\Lambda^{n}$ .

There are three important multiplicative bases for  $\Lambda^n$ .

Suppose that for each *n*,  $u_n$  is a symmetric function homogeneous of degree *n*. Then for any partition  $\lambda = (\lambda_1, \dots, \lambda_k)$ , we may define  $u_\lambda$  to be  $u_{\lambda_1} \cdots u_{\lambda_k}$ . There are three important multiplicative bases for  $\Lambda^n$ .

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If  $u_1, u_2, ...$  are algebraically independent, then  $\{u_{\lambda}\}_{\lambda \vdash n}$  will be a basis for  $\Lambda^n$ .

We define the *n*th elementary symmetric function  $e_n$  by

$$\boldsymbol{e}_n = \sum_{i_1 < \cdots < i_n} \boldsymbol{x}_{i_1} \cdots \boldsymbol{x}_{i_n},$$

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The *n*th complete symmetric function is

$$h_n=\sum_{i_1\leq\cdots\leq i_n}x_{i_1}\cdots x_{i_n},$$

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The *n*th power sum symmetric function is

$$p_n=\sum_{i=1}^{\infty}x_i^n,$$

so  $p_n = m_{(n)}$ .

Theorem. Each of  $\{h_{\lambda}\}_{\lambda \vdash n}$ ,  $\{e_{\lambda}\}_{\lambda \vdash n}$ , and  $\{p_{\lambda}\}_{\lambda \vdash n}$  is a basis for  $\Lambda^{n}$ .

# Some generating functions

We have

$$\sum_{n=0}^{\infty} e_n t^n = \prod_{i=1}^{\infty} (1+x_i t)$$

and

$$\sum_{n=0}^{\infty} h_n t^n = \prod_{i=1}^{\infty} (1 + x_i t + x_i^2 t^2 + \cdots)$$
$$= \prod_{i=1}^{\infty} \frac{1}{1 - x_i t}.$$

(Note that the variable *t* is redundant, since if we set t = 1 we can get it back by replacing each  $x_i$  with  $x_i t$ .)

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It follows that

$$\sum_{n=0}^{\infty} h_n t^n = \left(\sum_{n=0}^{\infty} (-1)^n e_n t^n\right)^{-1}.$$

Also

$$\log \prod_{i=1}^{\infty} \frac{1}{1 - x_i t} = \sum_{i=1}^{\infty} \log \frac{1}{1 - x_i t}$$
$$= \sum_{i=1}^{\infty} \sum_{n=1}^{\infty} x_i^n \frac{t^n}{n}$$
$$= \sum_{n=1}^{\infty} \frac{p_n}{n} t^n.$$

Therefore

$$\sum_{n=0}^{\infty} h_n t^n = \exp\left(\sum_{n=1}^{\infty} \frac{p_n}{n} t^n\right).$$

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Therefore

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If we expand the right side and equate coefficients of  $t^n$ 

then we get

$$h_n=\sum_{\lambda\vdash n}rac{p_\lambda}{z_\lambda}.$$

Here if  $\lambda = (1^{m_1} 2^{m_2} \cdots)$  then  $z_{\lambda} = 1^{m_1} m_1! 2^{m_2} m_2! \cdots$ .

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It is not hard to show that if  $\lambda$  is a partition of *n* then  $n!/z_{\lambda}$  is the number of permutations in the symmetric group  $\mathfrak{S}_n$  of cycle type  $\lambda$  and that  $z_{\lambda}$  is the number of permutations in  $\mathfrak{S}_n$  that commute with a given permutation of cycle type  $\lambda$ .

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For example, for n = 3 we have  $z_{(3)} = 3$ ,  $z_{(2,1)} = 2$ , and  $z_{(1,1,1)} = 6$ , so

$$h_3 = \frac{p_{(1,1,1)}}{6} + \frac{p_{(2,1)}}{2} + \frac{p_{(3)}}{3}$$
$$= \frac{p_1^3}{6} + \frac{p_2 p_1}{2} + \frac{p_3}{3}.$$

## The Cauchy kernel

The infinite product

$$\prod_{i,j=1}^{\infty} \frac{1}{1-x_i y_j}$$

is sometimes called the Cauchy kernel. It is symmetric in both  $x_1, x_2, \ldots$  and  $y_1, y_2, \ldots$ .

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First we note that the coefficient  $N_{\lambda,\mu}$  of  $x_1^{\lambda_1} x_2^{\lambda_2} \cdots y_1^{\mu_1} y_2^{\mu_2} \cdots$  in this product is the same as the coefficient of  $x_1^{\mu_1} x_2^{\mu_2} \cdots y_1^{\lambda_1} y_2^{\lambda_2} \cdots$ .

Now let's expand the product:

$$\prod_{i=1}^{\infty} \prod_{j=1}^{\infty} \frac{1}{1 - x_i y_j} = \prod_{i=1}^{\infty} \sum_{k=0}^{\infty} x_i^k h_k[y]$$

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Now  $N_{\lambda,\mu}$  is the coefficient of  $x_1^{\lambda_1} x_2^{\lambda_2} \cdots y_1^{\mu_1} y_2^{\mu_2} \cdots$  in this product, which is the same as the coefficient of  $y_1^{\mu_1} y_2^{\mu_2} \cdots$  in  $h_{\lambda}[y]$ .

# MacMahon's law of symmetry

Since  $N_{\lambda,\mu} = N_{\mu,\lambda}$ , we have MacMahon's law of symmetry: The coefficient of  $x_1^{\lambda_1} x_2^{\lambda_2} \cdots$  in  $h_{\mu}$  is equal to the coefficient of  $x_1^{\mu_1} x_2^{\mu_2} \cdots$  in  $h_{\lambda}$ .

# The scalar product

Now we define a scalar product on  $\Lambda$  by

$$\langle h_{\lambda}, f \rangle = \text{coefficient of } x_1^{\lambda_1} x_2^{\lambda_2} \cdots \text{ in } f$$

extended by linearity. By MacMahon's law of symmetry,  $\langle h_{\lambda}, h_{\mu} \rangle = \langle h_{\mu}, h_{\lambda} \rangle$ , so by linearity  $\langle f, g \rangle = \langle g, f \rangle$  for all  $f, g \in \Lambda$ .

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A short calculation shows that

$$\langle \boldsymbol{p}_{\lambda}, \boldsymbol{p}_{\mu} \rangle = \boldsymbol{z}_{\lambda} \delta_{\lambda,\mu}.$$

# The characteristic map

Let  $\rho$  be a representation of the symmetric group  $\mathfrak{S}_n$ ; i.e., an "action" of  $\mathfrak{S}_n$  on a finite-dimensional vector space *V* (over  $\mathbb{C}$ ).

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From  $\rho$  we can construct a function  $\chi^{\rho} : \mathfrak{S}_n \to \mathbb{C}$ , called the character of  $\rho$ , defined by

 $\chi^{
ho}(\boldsymbol{g}) = \operatorname{trace} \rho(\boldsymbol{g}).$ 

Then the character of  $\rho$  determines  $\rho$  up to equivalence.

We define the characteristic of  $\rho$  to be the symmetric function

$$\operatorname{ch} \rho = \frac{1}{n!} \sum_{g \in \mathfrak{S}_n} \chi^{\rho}(g) p_{\operatorname{cyc}(g)},$$

where cyc(g) is the cycle type of g.

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Since  $\chi^{\rho}(g)$  depends only on the cycle type of g, if we define  $\chi^{\rho}(\lambda)$ , for  $\lambda$  a partition of n, by  $\chi^{\rho}(\lambda) = \chi^{\rho}(g)$  for g with  $cyc(g) = \lambda$ , then we can write this as

$$\mathsf{ch} \, \rho = \frac{1}{n!} \sum_{\lambda \vdash n} \frac{n!}{z_{\lambda}} \chi^{\rho}(\lambda) p_{\lambda} \\ = \sum_{\lambda \vdash n} \chi^{\rho}(\lambda) \frac{p_{\lambda}}{z_{\lambda}}$$

Two very simple examples:

(1) The trivial representation. Here *V* is a one-dimensional vector space and for every  $g \in \mathfrak{S}_n$ ,  $\rho(g)$  is the identity transformation. Then  $\chi^{\rho}(g) = 1$  for all  $g \in \mathfrak{S}_n$  so

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(2) The regular representation. Here *V* is the vector space spanned by  $\mathfrak{S}_n$  and  $\mathfrak{S}_n$  acts by left multiplication. Then  $\chi^{\rho}(g) = n!$  if *g* is the identity element of  $\mathfrak{S}_n$  and  $\chi^{\rho}(g) = 0$  otherwise. So

ch 
$$ho = p_1^n$$
.

### Group actions

Let *G* be a finite group and let *S* be a finite set. An action of *G* on *S* is map  $\phi : G \times S \rightarrow S$ ,  $(g, s) \mapsto g \cdot s$  satisfying

•  $gh \cdot s = g \cdot (h \cdot s)$  for  $g, h \in G$  and  $s \in S$ 

•  $e \cdot s = s$  for all  $s \in S$ .

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Given an action of *G* on *S*, we get a representation of *G* on the vector space spanned by *S*:

$$\rho(g)\left(\sum_{s\in S}c_s\,s\right) = \sum_{s\in S}c_s\,g\cdot s$$

Then the trace of  $\rho(g)$  is the number of elements of *S* for which  $g \cdot s = s$ , which we denote by fix(*g*).

An important fact is Burnside's Lemma (also called the orbit-counting theorem): The number of orbits of *G* acting *S* is

$$\frac{1}{|G|}\sum_{g\in G} \mathsf{fix}(g).$$

Now we take *G* to be the symmetric group  $\mathfrak{S}_n$ .

The characteristic of the corresponding representation is

$$\frac{1}{n!}\sum_{g\in\mathfrak{S}_n}\operatorname{fix}(g)\,p_{\operatorname{cyc}(g)}=\sum_{\lambda\vdash n}\operatorname{fix}(\lambda)\,\frac{p_\lambda}{z_\lambda}$$

It is called the cycle index of the action of  $\mathfrak{S}_n$ , denoted  $Z_{\phi}$ .

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If we set all the  $p_i$  to 1 (or equivalently, set  $x_1 = 1$ ,  $x_i = 0$  for i > 0) then by Burnside's lemma we get the number of orbits. This is also equal to the scalar product  $\langle Z_{\phi}, h_n \rangle$ . There is a combinatorial interpretation to the coefficients of  $Z_{\phi}$ :

The coefficient of  $x_1^{\alpha_1} \cdots x_m^{\alpha_m}$  in  $Z_{\phi}$  is the number of orbits of the Young subgroup  $\mathfrak{S}_{\alpha} = \mathfrak{S}_{\alpha_1} \times \cdots \times \mathfrak{S}_{\alpha_m}$  of  $\mathfrak{S}_n$ , where  $\mathfrak{S}_{\alpha_1}$  permutes 1, 2, ...,  $\alpha_1$ ;  $\mathfrak{S}_{\alpha_2}$  permutates  $\alpha_1 + 1, \ldots, \alpha_1 + \alpha_2$ , and so on.

This coefficient is equal to the scalar product  $\langle Z_{\phi}, h_{\alpha} \rangle$ .

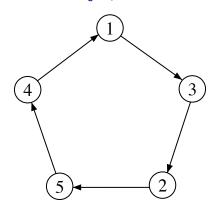
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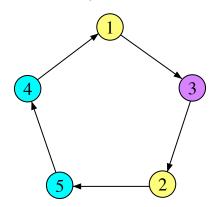
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This result is a form of Pólya's theorem. If  $\mathfrak{S}_n$  is acting on a set of "graphs" with vertex set  $\{1, 2, \ldots, n\}$  then we can construct the orbits of  $\mathfrak{S}_{\alpha}$  by coloring vertices  $1, 2, \ldots, \alpha_1$  in color 1; vertices  $\alpha_1 + 1, \ldots, \alpha_1 + \alpha_2$  in color 2, and so on, and then "erasing" the labels, leaving only the colors.

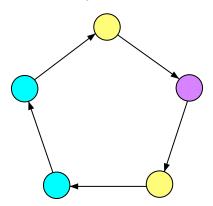
Example: The coefficient of  $x_1^2 x_2 x_3^2$  in the cycle index for  $\mathfrak{S}_5$  acting on directed 5-cycles,  $\frac{1}{5}(p_1^5 + 4p_5)$ 



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There are 6 of these.

# Schur functions

Another important basis for symmetric functions is the Schur function basis  $\{s_{\lambda}\}$ . The Schur functions are the characteristics of the irreducible representations of  $\mathfrak{S}_n$ , and they are orthonormal with respect to the the scalar product:

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They are, up to sign, the unique orthonormal basis that can be expressed as integer linear combinations of the  $m_{\lambda}$ .

If *f* is the characteristic of any representation of  $\mathfrak{S}_n$ , then *f* is a nonnegative integer linear combination of Schur functions.

One of them is called the Kronecker (or internal or inner) product. It has a very simple definition:

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It is extended by linearity to all symmetric functions, so

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The Kronecker product of symmetric functions corresponds to the tensor product of representations of  $\mathfrak{S}_n$ .

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First suppose that *g* can be expressed as a sum of monic terms, that is, monomials  $x_1^{\alpha_1} x_2^{\alpha_2} \dots$  with coefficient 1. For example,  $m_{\lambda}$  is a sum of monic terms.

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If we have a sum of monomials with positive integer coefficients then we can also write it as a sum of monic terms:

$$2p_2 = x_1^2 + x_1^2 + x_2^2 + x_2^2 + \cdots$$

In this case, if  $g = t_1 + t_2 + \cdots$ , where the  $t_i$  are monic terms, then

$$f[g] = f(t_1, t_2, ...).$$

For example

$$f[e_2] = f(x_1x_2, x_1x_3, x_2x_3, \dots)$$
  
$$f[2p_2] = f(x_1^2, x_1^2, x_2^2, x_2^2, \dots)$$

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More specifically,

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- For any g,  $p_n[g] = g[p_n]$
- $\blacktriangleright p_m[p_n] = p_{mn}$
- If *c* is a constant then  $c[p_n] = c$ .

These formulas allow us to define f[g] for any symmetric functions f and g.

## Examples of plethysm

First note that if c is a constant then

 $p_m[cp_n] = (cp_n)[p_m] = c[p_m]p_n[p_m] = cp_{mn}.$ 

Then since  $h_2 = (p_1^2 + p_2)/2$ , we have

$$h_2[-p_1] = \frac{1}{2}(p_1[-p_1]^2 + p_2[-p_1]) = \frac{1}{2}((-p_1)^2 - p_2) = e_2.$$

More generally, we can show that  $h_n[-p_1] = (-1)^n e_n$ . Also

$$h_2[1+p_1] = \frac{1}{2}(p_1[1+p_1]^2 + p_2[1+p_1])$$
  
=  $\frac{1}{2}((1+p_1)^2 + (1+p_2)) = 1 + p_1 + h_2$ 

Another example: Since

$$\prod_{i=1}^{\infty}(1+x_i)=\sum_{n=0}^{\infty}e_n,$$

we have

$$\prod_{i< j}(1+x_ix_j)=\sum_{n=0}^{\infty}e_n[e_2].$$

There are two special cases where we can often get simpler formulas for certain coefficients of symmetric functions, especially when they're expressed in terms of the power sums. There are two special cases where we can often get simpler formulas for certain coefficients of symmetric functions, especially when they're expressed in terms of the power sums.

First, the coefficient of  $x_1^n$  in a symmetric function f is the coefficient of  $x^n$  in f(x, 0, 0, 0), and if f is expressed in terms of the  $p_i$  we get this by setting  $p_i = x^i$  for all i.

Second, we can often get a simple formula or generating function for the coefficient of  $x_1 x_2 \cdots x_n$  in a symmetric function.

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Let E(f) be obtained from the symmetric function f (expressed in the  $p_i$ ) by setting  $p_1 = z$  and  $p_i = 0$  for all i > 1. Then

$$E(f)=\sum_{n=0}^{\infty}a_n\frac{z^n}{n!},$$

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Moreover, *E* is a homorphism,

E(f+g) = E(f) + E(g) and E(fg) = E(f)E(g),

and it respects composition,

$$E(f \circ g) = E(f) \circ E(g),$$

as long as g has no constant term.